

(2)

$$a_0 = \frac{1}{\sqrt{5}}, a_2 = \frac{2i}{\sqrt{5}} \text{ from a)}$$

$$\text{so } \Psi(x,t) = \frac{\psi_0(x)}{\sqrt{5}} \exp(-\frac{i\omega t}{2}) + \frac{2i\psi_2(x)}{\sqrt{5}} \exp(-\frac{5i\omega t}{2})$$

$$c) \langle E \rangle = \int_{-\infty}^{\infty} \Psi^*(x,t) \hat{H} \Psi(x,t) dx$$

This is unnecessarily complicated when there are only 2 terms.

$$= \int_{-\infty}^{\infty} \left( \sum_{n=0}^{\infty} a_n^* \psi_n^*(x) e^{+iE_n t/\hbar} \right) \hat{H} \left( \sum_{m=0}^{\infty} a_m \psi_m(x) e^{-iE_m t/\hbar} \right) dx$$

$$= \int_{-\infty}^{\infty} \left( \sum_{n=0}^{\infty} a_n^* \psi_n^*(x) e^{+iE_n t/\hbar} \right) \left( \sum_{m=0}^{\infty} a_m \hat{H} \psi_m(x) e^{-iE_m t/\hbar} \right) dx$$

$$= \int_{-\infty}^{\infty} \left( \sum_{n=0}^{\infty} a_n^* \psi_n^*(x) e^{+iE_n t/\hbar} \right) \left( \sum_{m=0}^{\infty} a_m E_m \psi_m(x) e^{-iE_m t/\hbar} \right) dx$$

(since the  $\psi_m$  are the eigenfunctions of the simple harmonic oscillator.)

We can write the above equation

$$\langle E \rangle = \int_{-\infty}^{\infty} \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} a_n^* a_m E_m \psi_n^*(x) \psi_m(x) e^{i(E_n - E_m)t/\hbar} dx$$

$$\langle E \rangle = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} a_n^* a_m E_m e^{i(E_n - E_m)t/\hbar} \int_{-\infty}^{\infty} \psi_n^*(x) \psi_m(x) dx$$

By orthonormality of eigenfunctions spatial parts

$$\int_{-\infty}^{\infty} \psi_n^*(x) \psi_m(x) dx = \begin{cases} 1, & m=n \\ 0, & m \neq n \end{cases}$$

$$\langle E \rangle = \sum_{n=0}^{\infty} a_n^* a_n E_n = \sum_{n=0}^{\infty} |a_n|^2 E_n$$

$$a_0 = \frac{1}{\sqrt{5}}, a_2 = \frac{2i}{\sqrt{5}}, a_n = 0 \text{ if } n \neq 0, 2$$

$$\therefore \langle E \rangle = \left| \frac{1}{\sqrt{5}} \right|^2 E_0 + \left| \frac{2i}{\sqrt{5}} \right|^2 E_2$$

the  $E_n$  are the harmonic oscillator energy eigenvalues  $E_n = (n + 1/2) \hbar \omega$ .

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