

for $R \leq r \leq \pi$, $V(r) = \frac{1}{4\pi\epsilon_0} \frac{2(Z-2)e^2}{r}$

$E = E_0$, the kinetic energy of the emitted α -particle

This equation is solved in FT. The solution is

$T \approx \exp\left(-\frac{\pi\sqrt{2m}}{h} \cdot \frac{V_0 R}{\sqrt{E}} + 4\sqrt{2mV_0} R\right)$

This is not relevant here.

which has the form

$T \approx A \exp(C/\sqrt{E})$ ①

This is all that is necessary.

There's a nicely compact relationship between the decay constant λ and the transmission coefficient T , given on FT page 402: $\lambda = \frac{v_1}{2R} T$, where v_1 is velocity of α -particle in nucleus ②

This says: probability of decay in 1 second equals the number of times the particle hits the potential wall at $r=R$ times the probability that it will pass through; a strong fuse of Newtonian and quantum mechanics.

Rearrange ② to give $T = \frac{2R\lambda}{v_1}$

① = ② $A \exp\left(\frac{C}{\sqrt{E}}\right) = \frac{2R\lambda}{v_1}$

This is the guide formula, but derivations have no place in an essay.

$\lambda = \frac{v_1 A}{2R} \exp\left(\frac{C}{\sqrt{E}}\right)$

$\ln \lambda = \frac{C}{\sqrt{E}} + \ln\left(\frac{v_1 A}{2R}\right)$

$\ln \lambda = -\frac{C}{\sqrt{E}} + K$ where $K = \ln\left(\frac{v_1 A}{2R}\right)$

This is a straight line graph of $\ln \lambda$ against $\frac{1}{\sqrt{E}}$, and observations

fit very well. If the line is plotted for different elements, with v_1 varying by 24 powers of 10, the slope of the line whose eqn. is given above does not differ from observations by more than a few per cent. I read in FT that α -particles do not exist inside the nucleus. I find this surprising. I would have thought

This is simple algebra in derivation on semi-classical.

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