

Consider the periodically forced system with Hamiltonian

$$H(q, p, t) = H_0(q, p) + \varepsilon q \sin \Omega t.$$

- (ii) Use part (i) and a result proved in *Unit 13* to deduce that, for this system, motion near the unperturbed separatrix is necessarily chaotic, provided that  $|\varepsilon|$  is sufficiently small.

[15]

#### Question 4

Consider the motion of a particle of mass  $m$  moving in the potential

$$V(q) = \begin{cases} 0, & |q| \leq L, \\ \infty, & |q| > L, \end{cases}$$

where  $L$  is a positive constant.

- (i) Show that all motion is bound and periodic, and that the action,  $I$ , the energy,  $E$ , and the frequency,  $\omega$ , are connected by the relations

$$E(I) = \frac{1}{2} A I^2, \quad \omega(I) = A I,$$

$$\text{where } A = \frac{\pi^2}{4mL^2}.$$

Show, further, that the angle variable,  $\theta$ , may be chosen so that for  $-\pi/2 \leq \theta \leq \pi/2$ ,

$$q = 2L\theta/\pi \quad \text{and} \quad q(\theta) = q(\pi - \theta). \quad [7]$$

- (ii) Show that the Fourier series for  $q(\theta, I)$  is

$$q(\theta, I) = \sum_{n=1}^{\infty} Q_n(I) \sin(2n+1)\theta,$$

$$\text{where } Q_n = \frac{8L}{\pi^2} \frac{(-1)^n}{(2n+1)^2}. \quad [6]$$

[Hint: You may find the following integral useful:

$$\int_0^{\pi/2} d\theta \theta \sin(2n+1)\theta = \frac{1}{(2n+1)^2} \cos n\pi.]$$

- (iii) The system is perturbed by a periodic force, so that the Hamiltonian becomes

$$H(q, p, t) = \frac{p^2}{2m} + V(q) + Fq \sin \Omega t.$$

Show that the  $n$ th resonance is at the action

$$I_n = \frac{4m\Omega L^2}{\pi^2(2n+1)},$$

and that near here the Hamiltonian can be cast into the form

$$K(\phi, J) = \frac{1}{2} A J^2 + \frac{4LF}{\pi^2(2n+1)^2} \sin(2n+1)\phi.$$

Hence, or otherwise, deduce that the separatrix width  $S^{(n)}$  of the  $n$ th resonance island is given by

$$S^{(n)} = \frac{16\sqrt{mFL^3}}{(2n+1)\pi^2}. \quad [8]$$

- (iv) Using the result of *Unit 11* Section 11.5, that for resonance islands not to intersect we must have  $\Delta I_n > \frac{1}{2}(S^{(n)} + S^{(n+1)})$ , deduce that for large  $n$  the resonant approximation is valid near the  $n$ th resonance only if

$$F < \frac{mL\Omega^2}{16n^2}. \quad [4]$$